

Forcing of Solutions to Reaction-Diffusion Equations With Applications to Population Models

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Introduction

Initial-boundary value problem:

$$(1) \quad \mathbf{u}_t = \mathbf{f}(\mathbf{u}) + \mathbf{g}(\mathbf{x}, t) + D\Delta\mathbf{u} \quad \text{on } \Omega \times (0, \infty)$$

$$(2) \quad \frac{\partial \mathbf{u}}{\partial n}(\mathbf{x}, t) = 0 \quad \text{on } \partial\Omega \times (0, \infty)$$

$$(3) \quad \mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0(\mathbf{x}) \quad \text{on } \Omega$$

- Reaction-diffusion equation's solution converges to a spatially homogeneous solution
- Analysis of the asymptotic behavior simplified

Goal

- Find solution to a system of ODEs with spatially averaged reaction terms.
- Show this solution is stable as a solution to the reaction-diffusion equations.
- Show that when the diffusion coefficients are sufficiently large, solutions are asymptotically homogeneous in Ω .

Spatially Averaged Equation

Solution to (1)–(3) will be compared to the solution of:

$$(4) \quad \mathbf{v}_t = \mathbf{f}(\mathbf{v}) + \bar{\mathbf{g}}(t) \quad \text{for } t > 0$$

$$(5) \quad \mathbf{v}(0) = \frac{1}{|\Omega|} \int_{\Omega} \mathbf{u}_0(\mathbf{x}) d\Omega$$

where

$$\bar{\mathbf{g}}(t) = \frac{1}{|\Omega|} \int_{\Omega} \mathbf{g}(\mathbf{x}, t) d\Omega$$

and $|\Omega|$ is the measure of the spatial domain Ω .

Stability

Definition 1 *A bounded solution $\mathbf{v}(t)$, to (4) and (5) will be said to be stable as a solution to (1) and (2) if for every $\epsilon > 0$ there exists a $\delta > 0$ and a time $T > 0$ such that*

$$\|\mathbf{u}(\mathbf{x}, 0) - \mathbf{v}(0)\|_{L^2(\Omega)} < \delta$$

implies

$$\|\mathbf{u}(\mathbf{x}, t) - \mathbf{v}(t)\|_{L^2(\Omega)} \leq \epsilon$$

for all $t \geq T$, where

$$\|\cdot\|_{L^2(\Omega)}^2 = \int_{\Omega} \langle \cdot, \cdot \rangle d\Omega$$

Stability Result

Theorem 2 *Suppose there exists a compact positively invariant region $\Sigma \subset \mathbb{R}^n$ for (1) with $u_0(x) \in \Sigma$ and suppose $v(t)$ is a solution to (4) such that there exists a positive constant γ for which $\langle f_u(v(t))w, w \rangle \leq -\gamma|w|^2$ for all $w \in \Sigma$, then if d , the smallest diffusion coefficient of D is sufficiently large, $v(t)$ is stable as a solution to (1) and (2).*

Outline of Proof

1. Bound the deviation of the solution to (1)–(3) from its spatial average over domain Ω .
2. Bound the deviation of the spatial average from the solution to (4) and (5).
3. Use the Minkowski inequality to combine those two bounds into a bound on the deviation of the solution to (1)–(3) from that of (4) and (5).

Spatially Averaged Solution

Spatial average of solution to (1)–(3) is

$$\bar{\mathbf{u}}(t) = |\Omega|^{-1} \int_{\Omega} \mathbf{u}(\mathbf{x}, t) d\Omega,$$

Solves the following initial value problem for $t > 0$.

$$(6) \quad \bar{\mathbf{u}}_t = \mathbf{f}(\bar{\mathbf{u}}) + \bar{\mathbf{g}}(t) + \frac{1}{|\Omega|} \int_{\Omega} \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}) d\Omega$$

$$(7) \quad \bar{\mathbf{u}}(0) = \frac{1}{|\Omega|} \int_{\Omega} \mathbf{u}_0(\mathbf{x}) d\Omega$$

Energy Method

Let $\phi(t) = (1/2)\|\mathbf{u}(\mathbf{x}, t) - \bar{\mathbf{u}}(t)\|_{L^2(\Omega)}^2$, then

$$\begin{aligned}\frac{d\phi}{dt} &= \int_{\Omega} \langle \mathbf{u}_t - \bar{\mathbf{u}}_t, \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega \\ &= \int_{\Omega} \langle \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}), \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega + \\ &\quad \int_{\Omega} \langle \mathbf{g}(\mathbf{x}, t) - \bar{\mathbf{g}}(t), \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega + \\ &\quad \int_{\Omega} \left\langle \frac{1}{|\Omega|} \int_{\Omega} \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}) d\Omega, \mathbf{u} - \bar{\mathbf{u}} \right\rangle d\Omega + \\ &\quad \int_{\Omega} \langle D\Delta\mathbf{u}, \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega\end{aligned}$$

Diffusion Integral

$$\begin{aligned}\int_{\Omega} \langle D\Delta \mathbf{u}, \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega &= - \int_{\Omega} \langle D\nabla(\mathbf{u} - \bar{\mathbf{u}}), \nabla(\mathbf{u} - \bar{\mathbf{u}}) \rangle d\Omega \\ &\leq -d \|\nabla(\mathbf{u} - \bar{\mathbf{u}})\|_{L^2(\Omega)}^2\end{aligned}$$

Since $\bar{\mathbf{u}}$ is independent of \mathbf{x} , by Green's first identity, and the boundary conditions of equation (2). Then

$$(8) \quad \int_{\Omega} \langle D\Delta \mathbf{u}, \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega \leq -d\lambda \|\mathbf{u} - \bar{\mathbf{u}}\|_{L^2(\Omega)}^2 = -2d\lambda\phi(t),$$

λ is the smallest positive eigenvalue of $-\Delta$ with homogeneous Neumann boundary conditions on Ω .

Reaction Term Integral

$$\begin{aligned}\langle \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}), \mathbf{u} - \bar{\mathbf{u}} \rangle &\leq \frac{1}{2} |\mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}})|^2 + \frac{1}{2} |\mathbf{u} - \bar{\mathbf{u}}|^2 \\ &\leq \frac{M^2}{2} |\mathbf{u} - \bar{\mathbf{u}}|^2 + \frac{1}{2} |\mathbf{u} - \bar{\mathbf{u}}|^2,\end{aligned}$$

where $M = \max_{\mathbf{u} \in \Sigma} \{|\mathbf{f}_{\mathbf{u}}|\}$. Thus the integral below involving the reaction term of (1) obeys the inequality:

$$(9) \quad \int_{\Omega} \langle \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}), \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega \leq (M^2 + 1)\phi(t).$$

Forcing Term Integral

$$\langle \mathbf{g}(\mathbf{x}, t) - \bar{\mathbf{g}}(t), \mathbf{u} - \bar{\mathbf{u}} \rangle \leq \frac{1}{2} |\mathbf{g}(\mathbf{x}, t) - \bar{\mathbf{g}}(t)|^2 + \frac{1}{2} |\mathbf{u} - \bar{\mathbf{u}}|^2$$

Components of \mathbf{g} and $\bar{\mathbf{g}}$ are periodic in t , and \mathbf{g} is continuous in \mathbf{x} , thus there exists $K > 0$ such that

$$\|\mathbf{g}(\mathbf{x}, t) - \bar{\mathbf{g}}(t)\|_{L^2(\Omega)}^2 \leq K^2.$$

Therefore

$$(10) \quad \int_{\Omega} \langle \mathbf{g}(\mathbf{x}, t) - \bar{\mathbf{g}}(t), \mathbf{u} - \bar{\mathbf{u}} \rangle d\Omega \leq \frac{K^2}{2} + \phi(t)$$

Averaged Integral

$$(11) \quad \int_{\Omega} \left\langle \int_{\Omega} \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}) d\Omega, \mathbf{u} - \bar{\mathbf{u}} \right\rangle d\Omega = \\ \left\langle \int_{\Omega} \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}) d\Omega, \int_{\Omega} \mathbf{u} - \bar{\mathbf{u}} d\Omega \right\rangle = 0$$

Finally we have

$$(12) \quad \frac{d\phi}{dt} + (2d\lambda - M^2 - 2)\phi(t) \leq \frac{K^2}{2}$$

Let $\sigma = 2d\lambda - M^2 - 2$.

Differential Inequality

Integrating (12) with respect to time over the interval $[0, t]$ produces

$$\phi(t) \leq \frac{1}{2} \|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 e^{-\sigma t} + \frac{K^2}{2\sigma} (1 - e^{-\sigma t})$$

or equivalently

$$(13) \quad \|\mathbf{u}(\mathbf{x}, t) - \bar{\mathbf{u}}(t)\|_{L^2(\Omega)}^2 \leq \|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 e^{-\sigma t} + \frac{K^2}{\sigma}.$$

2nd Step

Define $\psi(t) = (1/2)|\bar{\mathbf{u}}(t) - \mathbf{v}(t)|^2$.

$$\begin{aligned}\frac{d\psi}{dt} &= \langle \bar{\mathbf{u}}_t - \mathbf{v}_t, \bar{\mathbf{u}} - \mathbf{v} \rangle \\ &\leq -\gamma|\bar{\mathbf{u}} - \mathbf{v}|^2 + \frac{1}{|\Omega|} \left\langle \int_{\Omega} \mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}}) d\Omega, \bar{\mathbf{u}} - \mathbf{v} \right\rangle \\ &\quad + \langle \mathbf{f}(\bar{\mathbf{u}}) - \mathbf{f}(\mathbf{v}) - \mathbf{f}_{\mathbf{u}}(\mathbf{v})(\bar{\mathbf{u}} - \mathbf{v}), \bar{\mathbf{u}} - \mathbf{v} \rangle\end{aligned}$$

For every $\delta > 0$ there exists an $\eta \equiv \eta(\delta) > 0$ such that if $|\bar{\mathbf{u}} - \mathbf{v}|^2 < 2\eta$ then

$$|\langle \mathbf{f}(\bar{\mathbf{u}}) - \mathbf{f}(\mathbf{v}) - \mathbf{f}_{\mathbf{u}}(\mathbf{v})(\bar{\mathbf{u}} - \mathbf{v}), \bar{\mathbf{u}} - \mathbf{v} \rangle| < \frac{\delta}{2}|\bar{\mathbf{u}} - \mathbf{v}|^2.$$

Note $\psi(0) = 0$ and $\psi(t)$ is continuous. There exists $t_0 > 0$ for which $\psi(t) < \eta$ on the interval $[0, t_0)$. Thus

$$\frac{d\psi}{dt} \leq \left(\delta - 2\gamma + \frac{\delta}{|\Omega|} \right) \psi(t) + \frac{1}{2\delta|\Omega|} \left(\int_{\Omega} |\mathbf{f}(\mathbf{u}) - \mathbf{f}(\bar{\mathbf{u}})| d\Omega \right)^2.$$

Choose $\delta = \gamma|\Omega|/(|\Omega| + 1)$, use Jensen's inequality, and the inequality in (13), then on the interval $[0, t_0)$,

$$(14) \quad \frac{d\psi}{dt} + \gamma\psi(t) \leq \frac{M^2(|\Omega| + 1)}{2\gamma|\Omega|} \left(\|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 e^{-\sigma t} + \frac{K^2}{\sigma} \right)$$

Let $N^2 = (M^2/\gamma)(1 + |\Omega|^{-1})$. Then for $0 \leq t < t_0$,

$$(15) \quad \psi(t) \leq \frac{N^2}{2(\gamma - \sigma)} \|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 (e^{-\sigma t} - e^{-\gamma t}) + \frac{K^2 N^2}{2\gamma\sigma} (1 - e^{-\gamma t}).$$

Choose $d > (\gamma + M^2 + 2)/(2\lambda)$ then $\sigma > \gamma > 0$ and $0 < e^{-(\sigma-\gamma)t} \leq 1$ for all $t \geq 0$. Thus for $t \in [0, t_0)$

$$(16) \quad |\bar{\mathbf{u}}(t) - \mathbf{v}(t)|^2 \leq \frac{N^2}{\sigma - \gamma} \|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 e^{-\gamma t} + \frac{K^2 N^2}{\gamma\sigma}.$$

Separation Norm

By a first-time argument the previous inequality holds for all $t \geq 0$. Hence

$$\|\bar{\mathbf{u}}(t) - \mathbf{v}(t)\|_{L^2(\Omega)}^2 \leq \frac{N^2|\Omega|}{\sigma - \gamma} \|\mathbf{u}_0(\mathbf{x}) - \bar{\mathbf{u}}(0)\|_{L^2(\Omega)}^2 e^{-\gamma t} + \frac{K^2 N^2 |\Omega|}{\gamma \sigma}$$

Separation of solutions (1)–(3) from the solutions of (4) and (5) is governed by

$$\|\mathbf{u}(\mathbf{x}, t) - \mathbf{v}(t)\|_{L^2(\Omega)} \leq A_1 \|\mathbf{u}_0(\mathbf{x}) - \mathbf{v}(0)\|_{L^2(\Omega)} e^{-\gamma t/2} + A_2$$

where A_1 and A_2 are positive constants.

Effects of Forcing

PDE:

$$(17) \quad \mathbf{u}_t = \mathbf{f}(\mathbf{u}) + \mathbf{g}(\mathbf{x}, t; \mathbf{p}) + D\Delta\mathbf{u} \quad \text{on } \Omega \times (0, \infty)$$

$$(18) \quad \mathbf{u}(\mathbf{x}, 0) = \mathbf{u}_0(\mathbf{x}) \quad \text{on } \Omega$$

Spatial average:

$$(19) \quad \mathbf{v}_t = \mathbf{f}(\mathbf{v}) + \bar{\mathbf{g}}(t; \mathbf{p})$$

$$(20) \quad \mathbf{v}(0) = \frac{1}{|\Omega|} \int_{\Omega} \mathbf{u}_0(\mathbf{x}) d\Omega$$

Temporal average:

$$(21) \quad \mathbf{w}_t = \mathbf{f}(\mathbf{w}) + \mathbf{A}$$

$$(22) \quad \mathbf{w}(0) = \mathbf{v}(0)$$

Separation Result for Equilibria

Theorem 3 *If $f \in C^1(\Omega)$ where $\Omega \subset \mathbb{R}^n$ and e is a stable hyperbolic equilibrium of (21), and if $\bar{g} = \langle \bar{g}_1, \dots, \bar{g}_n \rangle$ where \bar{g}_i is continuous and periodic in t with period p_i and has time average over a period of A_i for $i = 1, \dots, n$, and if there exists a compact positively invariant region $\Sigma \subset \mathbb{R}^n$ for (17) with $u_0(x) \in \Sigma$ and if $v(t)$ is a solution to (19) such that there exists a positive constant γ for which $\langle f_u(v(t))z, z \rangle \leq -\gamma|z|^2$ for all $z \in \Sigma$, then when the smallest diffusion coefficient of D is sufficiently large, e is stable as a solution to (17) with homogeneous Neumann boundary conditions.*

Separation Norm

$$\begin{aligned} \|\mathbf{u}(\mathbf{x}, t) - \mathbf{e}\|_{L^2(\Omega)} &\leq A_1 \|\mathbf{u}_0(\mathbf{x}) - \mathbf{v}(0)\|_{L^2(\Omega)} e^{-\gamma t/2} + A_2 + \\ &\quad B_1 |\Omega|^{1/2} \sum_{i=1}^n p_i \|\bar{\mathbf{g}}_i\|_{\infty} + \\ &\quad B_2 |\Omega|^{1/2} |\mathbf{v}(0) - \mathbf{e}| e^{-\alpha t} \end{aligned}$$

Separation Result for Periodic Orbits

Theorem 4 *Suppose $f \in C^1(\Omega)$ where $\Omega \subset \mathbb{R}^n$ and $w(t)$ is an asymptotically stable periodic solution of (21). Let the forcing term $\bar{g} = \langle \bar{g}_1, \dots, \bar{g}_n \rangle$ where \bar{g}_i is continuous and periodic in t with period p_i and time average over a period of A_i for $i = 1, \dots, n$. Suppose there exists a compact positively invariant region $\Sigma \subset \mathbb{R}^n$ for (17) with $u_0(x) \in \Sigma$ and suppose $v(t)$ is a solution to (19) such that there exists a positive constant γ for which $\langle f_u(v(t))z, z \rangle \leq -\gamma|z|^2$ for all $z \in \Sigma$. Then when the smallest diffusion coefficient of the diagonal diffusion matrix is sufficiently large, $w(t)$ is stable as a solution to (17) with homogeneous Neumann boundary conditions.*

Example

$$u_t = u(1 - c_{11}u - v) + g_1(x, t) + D_1u_{xx}$$

$$v_t = v(0.5 - u - c_{22}v)(u + c_{22}v - 1.5) + g_2(x, t) + D_2v_{xx}$$

Zero-flux boundary conditions are assumed.

- $c_{11} = 1.1, c_{22} = 0.35,$

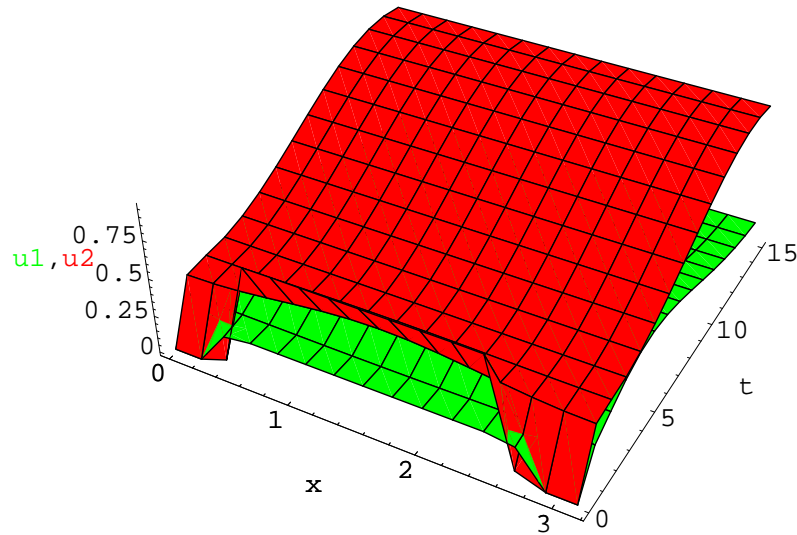
- $D_1 = 1, D_2 = 1/2,$

- and $g_1 = g_2 = 0.$

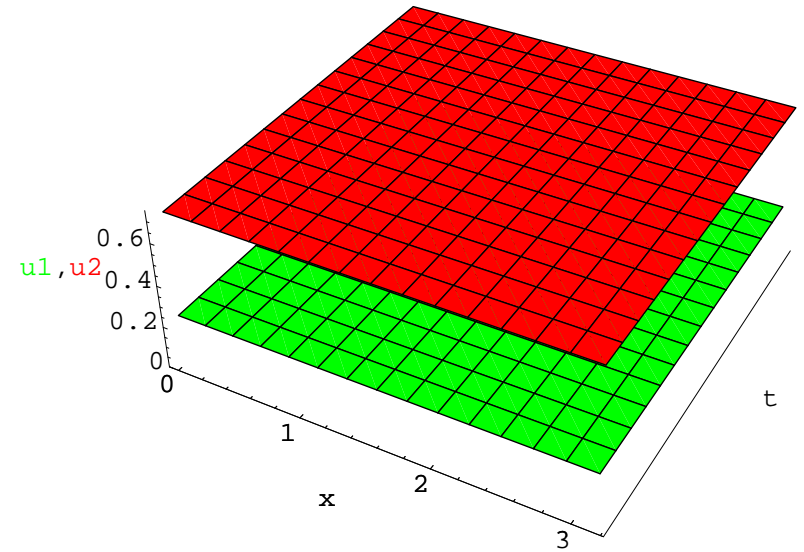
Asymptotically stable equilibrium solution

$$e \approx (0.2439, 0.7317).$$

Graphs



(a) $0 \leq t \leq 15$



(b) As $t \rightarrow \infty$

The solution to a set of reaction-diffusion equations exponentially approaches a constant equilibrium solution.

Hopf Bifurcation

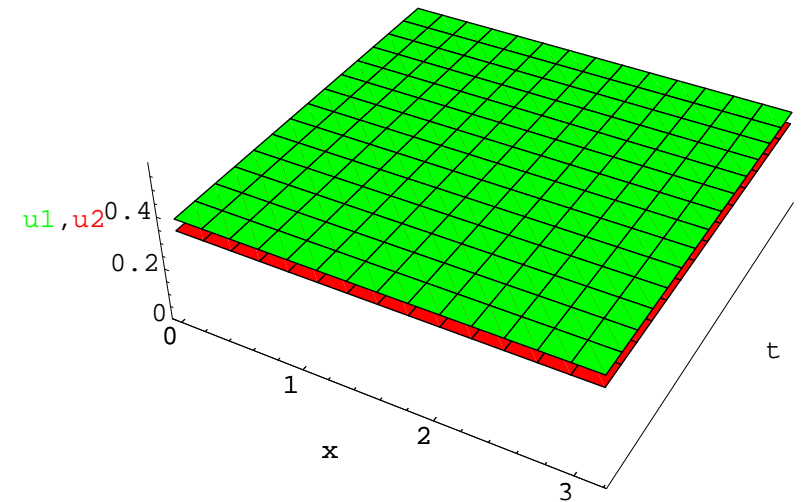
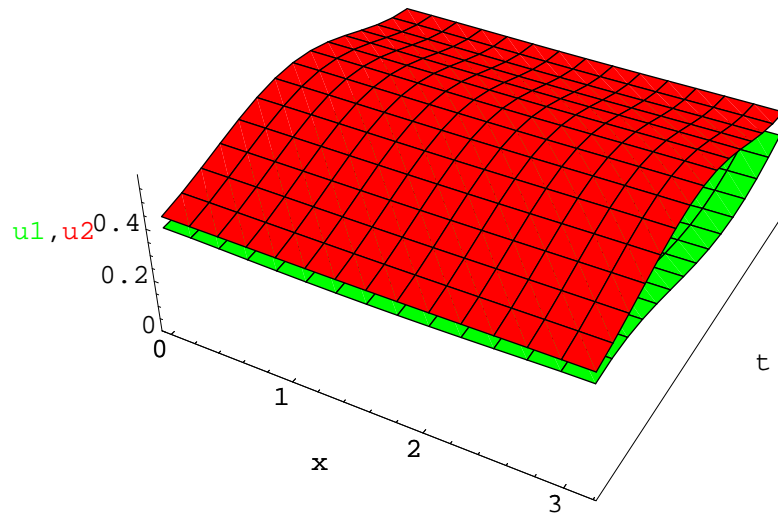
As c_{11} is decreased to 1.05, the equilibrium solution undergoes a Hopf bifurcation. Nearby an attracting periodic orbit appears.

Harvesting u at the constant rate of -0.1 over all of Ω will restabilize an equilibrium at $e \approx (0.3820, 0.3371)$.

The system can be brought back near the equilibrium point with time-periodic harvesting that is nonuniformly distributed in space and whose average over Ω and a period is -0.1 .

Let $g_1(x, t) = 0.6580B(x)(-0.1 + 0.05 \cos 2\pi t)$ where $B(x)$ is a $C^\infty(\mathbb{R})$ “bump function”.

Graphs



(a) Spatially varying, periodic forcing (b) Constant forcing

Comparison of the asymptotic separation of solutions of system of reaction-diffusion equations subject to varying and constant forcing.